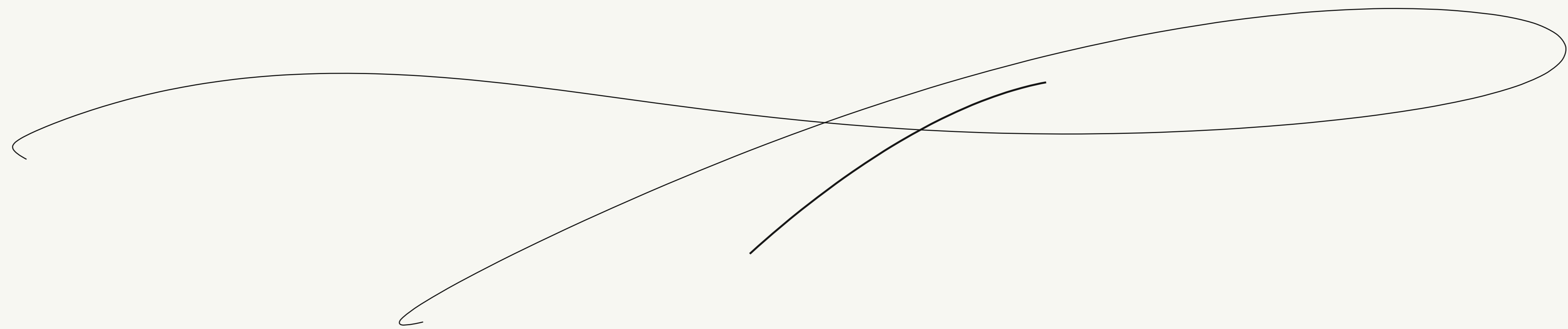


Towards a Geometry and Analysis
for Bayesian Mechanics



07 March 2022

Dutton A R Saktiwindri

Formal structures in physics:

↳ Geometry - study of spaces and their properties,
such as the state spaces of dynamical systems

Analysis - study of functions and their
properties, such as the equations of motion of
dynamical systems

Motivation: relate ^{analysis} gauge symmetries and ^{geometric} gauge forces to inference and complex systems, especially in the context of the free energy principle Bayesian mech.

Approach:

- ① Relate the FEP to maximum entropy
- ② Show max ent is actually a gauge theory
- ③ Profit?

Overview of some recent results

slides can be found later at darsakthi.github.io/talks; preprints to follow (soon)

Find p arg max $\left(- \int \ln p p \right)$

subject to $\int \int_{\mathbb{R}} f(x) p(x) dx = c$

$\hookrightarrow \mathbb{E}[f(X)] = c$

$$\text{Find } p \quad - \int \ln p p - \lambda \left(\int y p + C \right)$$

is max

$$\frac{\partial}{\partial p} \left(- \ln p p - \lambda y p \right) = 0$$

~~$$- \ln p - \lambda y = 0$$~~

~~$$p = \exp \{ - \lambda y \}$$~~

$$\lambda y - \lambda y = 0 \quad \checkmark$$

$$\ell - \beta V(x)$$

$V(x)$: a potential

$$V(x) = \int f(x)$$

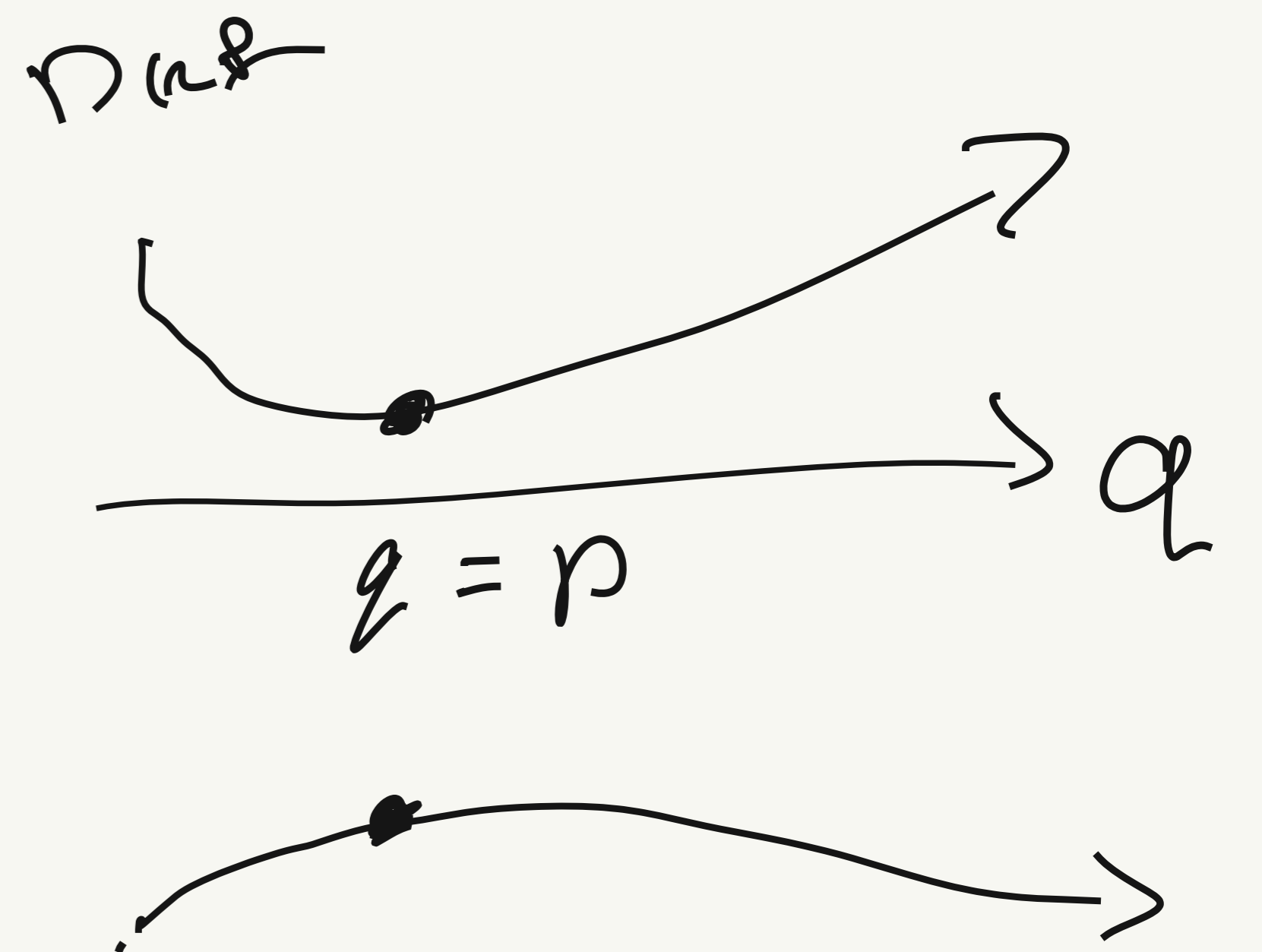
$$\left. \begin{aligned} & \partial_x (\partial_x \psi) \\ & - \beta \partial_{xx} \rho \\ & = 0 \end{aligned} \right\}$$

$$\hookrightarrow D_{KL}(q||p) = \mathbb{E}_q[\ln q] - \mathbb{E}_q[\ln p]$$

$$\arg \min_q (D_{KL}) \Rightarrow q = p$$

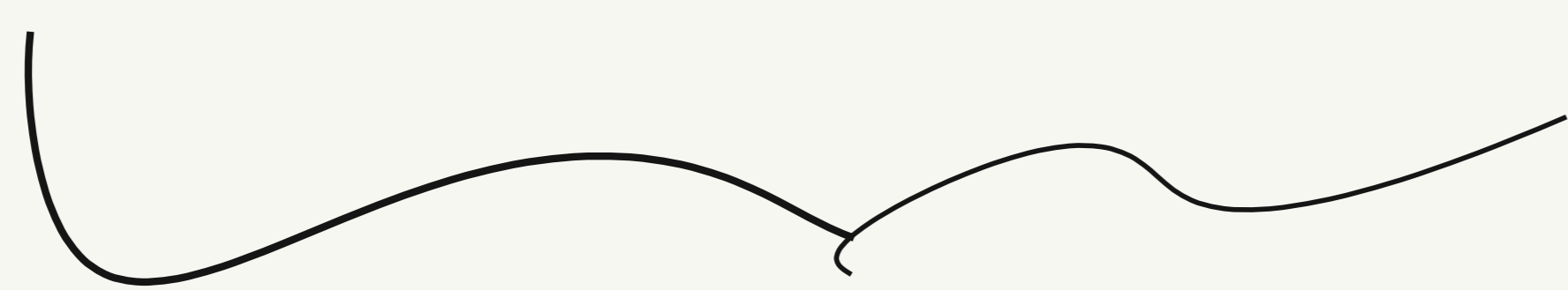
$$\hookrightarrow \arg \max_q (-D_{KL}) \Rightarrow q = p$$

$$-\mathbb{E}_q \ln q + \mathbb{E}_q \ln p$$



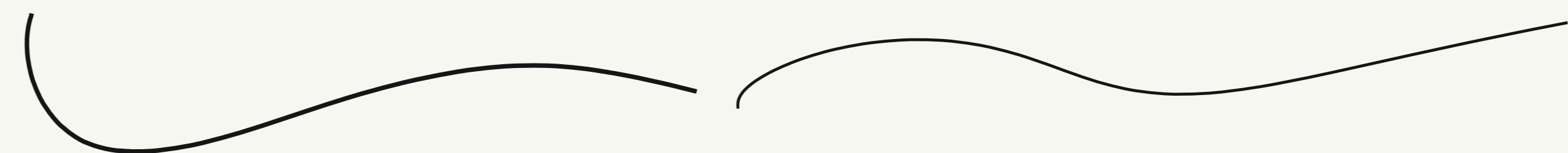
$$-\int \ln q \, q + \int \ln p \, q = 0$$

$$-\int \ln q \, q - \int -\ln p \, q + 0$$



entropy

\mathcal{L}



$$E[-\ln p] = 0$$

$$-\int \ln q \, q - \left(\int -\ln p(n|\mu, s) \, q + \ln p(\mu, s) \right)$$

$$\mathbb{E} [-\ln p(n|\mu, s)]$$

$$= -\ln p(\mu, s)$$

Thm. For an appropriate set of constraints,

$$FEP \iff \text{Max Ent.}$$

Proof. $\arg \max_q (-D_{KL}) = \arg \min_q (D_{KL})$

$$q^* = p$$

$$q^* = Z \exp \left\{ - \left(-\ln p \right) \right\}$$

$q^* = p$

Symmetry of Markov blanket

a
hidden
assumption

$$\begin{cases} \sigma(\bar{\mu}) = \bar{\pi} \\ \varrho^{-1}(\bar{\pi}) = \bar{\mu} \end{cases}$$

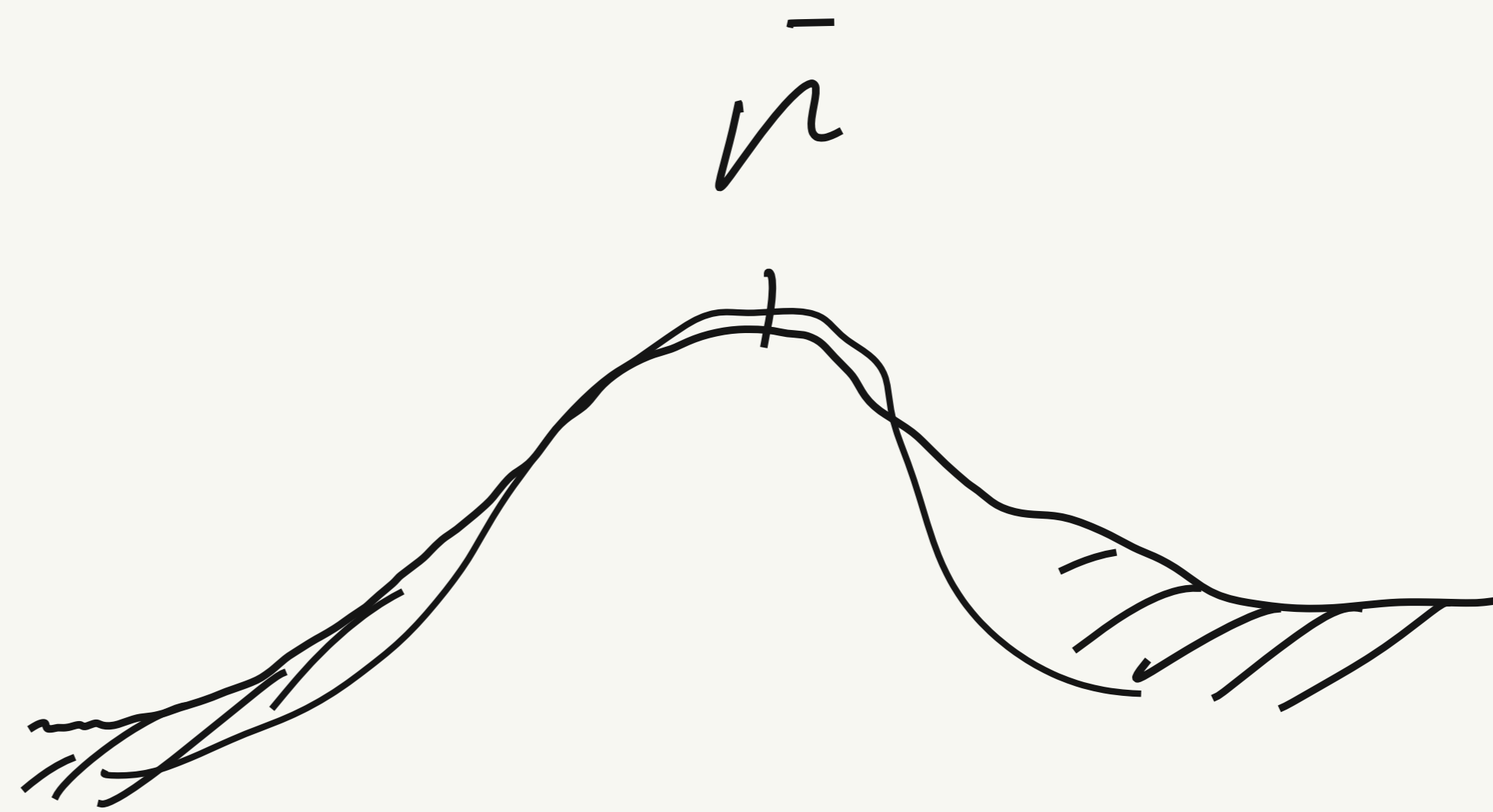
If $q = q(\mu; \sigma(\bar{\mu}))$ then q minimizes D_{KL} .

Proof. Suppose p has $\bar{\mu}$ as a suff. stat.

then $p = \sum e^{r_i}$ and $\mathbb{E}_p[r] = \bar{\mu}$.

Since $p = p(\mu; \bar{\mu})$ and $\sigma(\bar{\mu})$

$= \bar{\mu}$ $q = p$



$q \approx p$

$J = \mu$ and

$$E[\mu] = \sigma(\bar{\mu})$$

$\mu \mapsto \mu$

$$\Rightarrow E \mu = \sigma^{-1}(\bar{\mu})$$

max (S) Subject to

μ model μ

or neg.

Introducing Constraints decreases entropy

BUT does not minimize it.

$$J = x^2 + y^2$$

2D Gaussian

$p \propto -J$



$$J = \text{Constraint}$$

$P = \text{prob.}$

Indeed

$$J = x^2 + y^2$$

$$\|\vec{x}\|$$

$$\mu = [x, y]$$

$$g(\mu) = \exp\{-\lambda J\}$$

$$l = -x^2 - y^2 = \mathcal{N}(0, \lambda)$$

$p(\text{state}) \approx$ — constraints

Via Max entropy / min FE

Constraints on states of a system are like preferences about what states are likely to be occupied



which are like potential functions for sampling dynamics.

So what?

Potential functions are geometric

features of a system

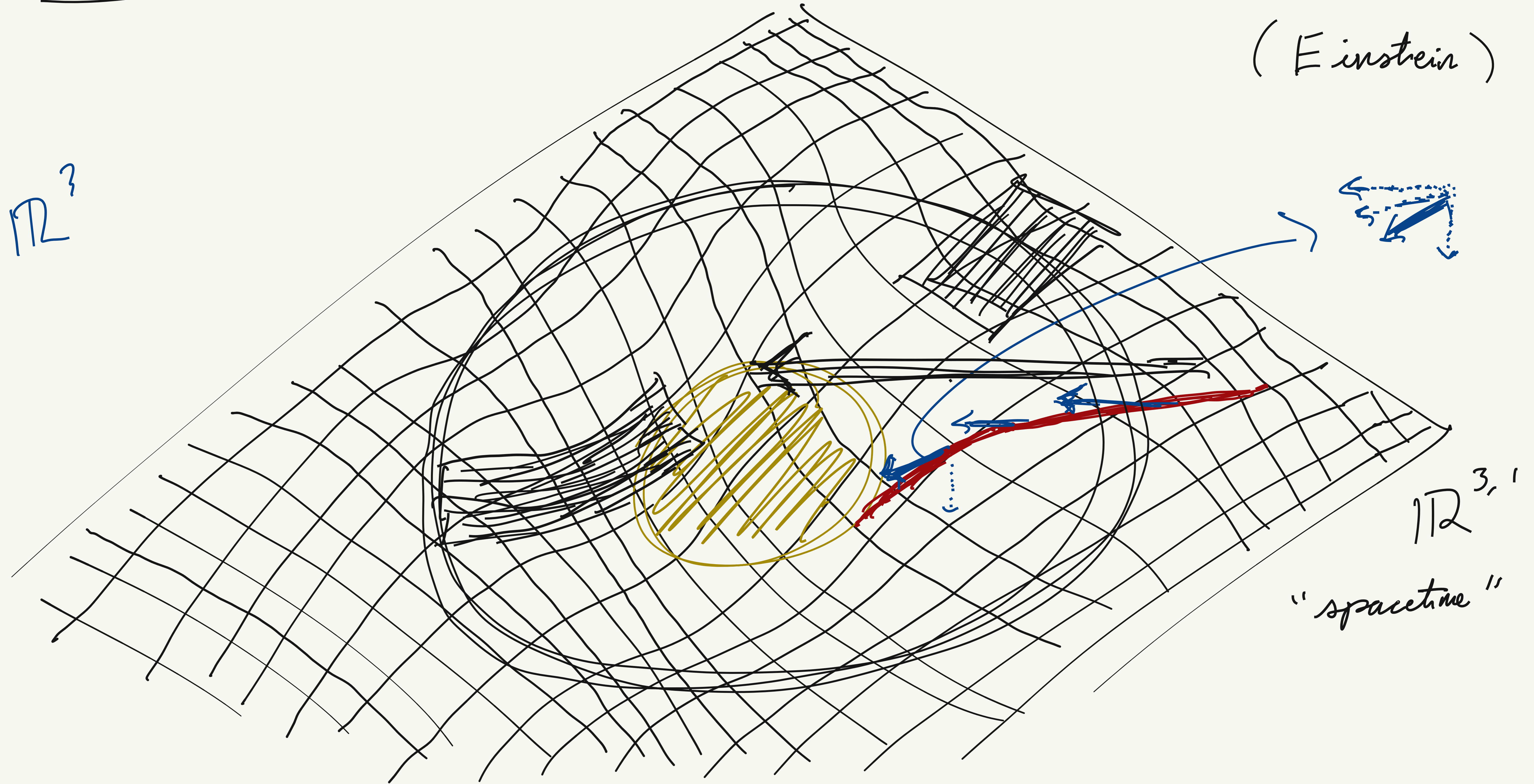
example: general relativity

Some geometry

Equivalence principle

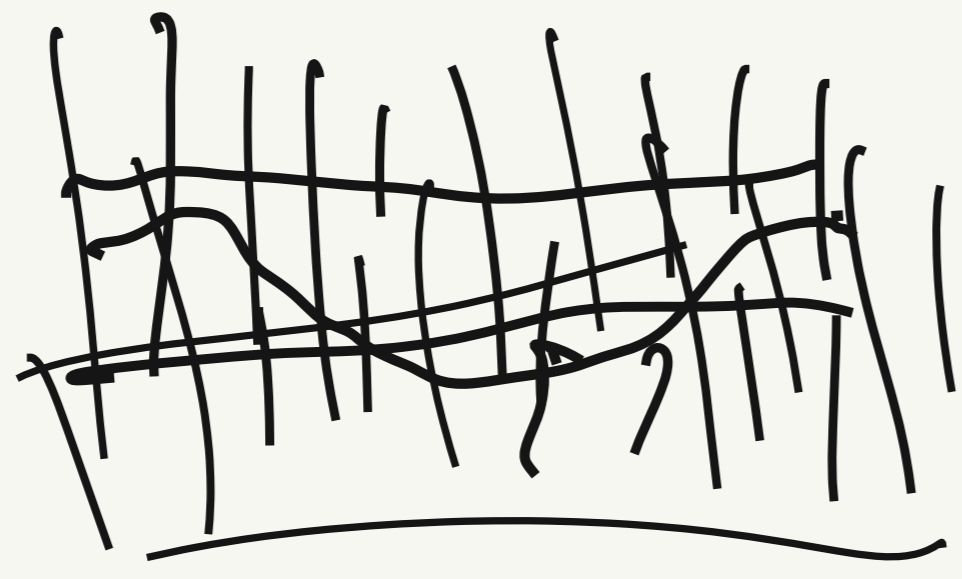
(Einstein)

\mathbb{R}^3

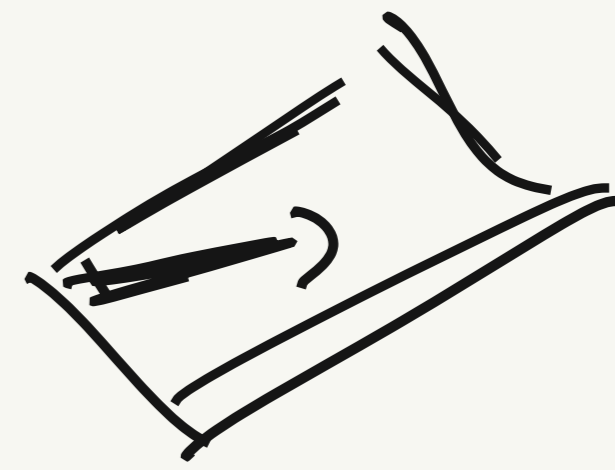
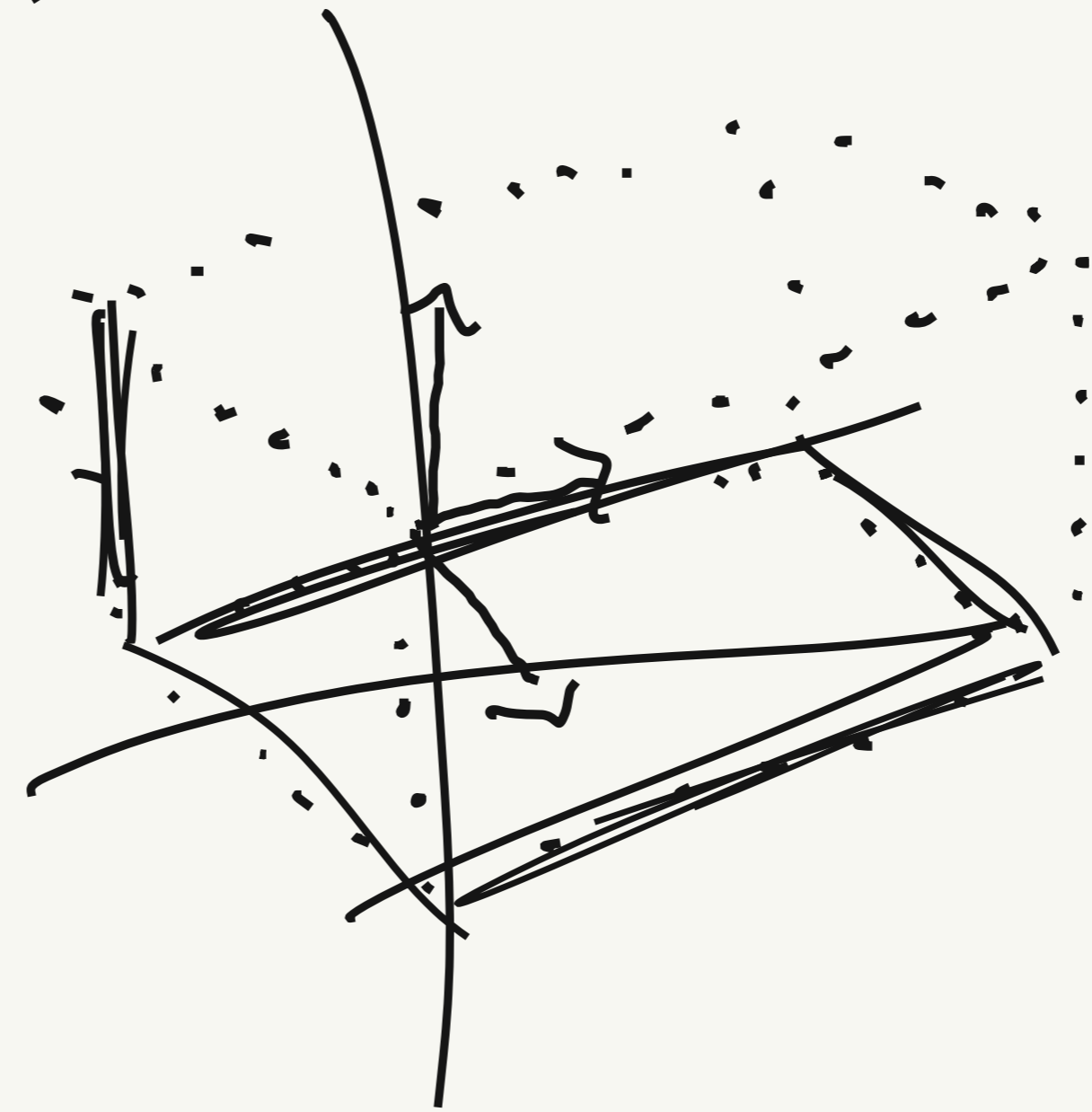


$\mathbb{R}^{3,1}$

"spacetime"



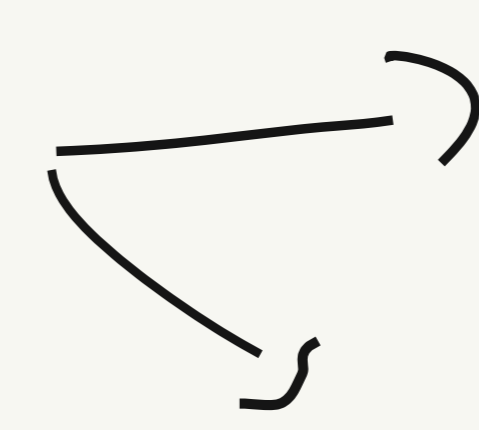
$$\mathbb{R}^2 + \mathbb{R} = \mathbb{R}^3$$



$$H_x T$$



$$V_x T$$



Spacetime



force
(gravity)

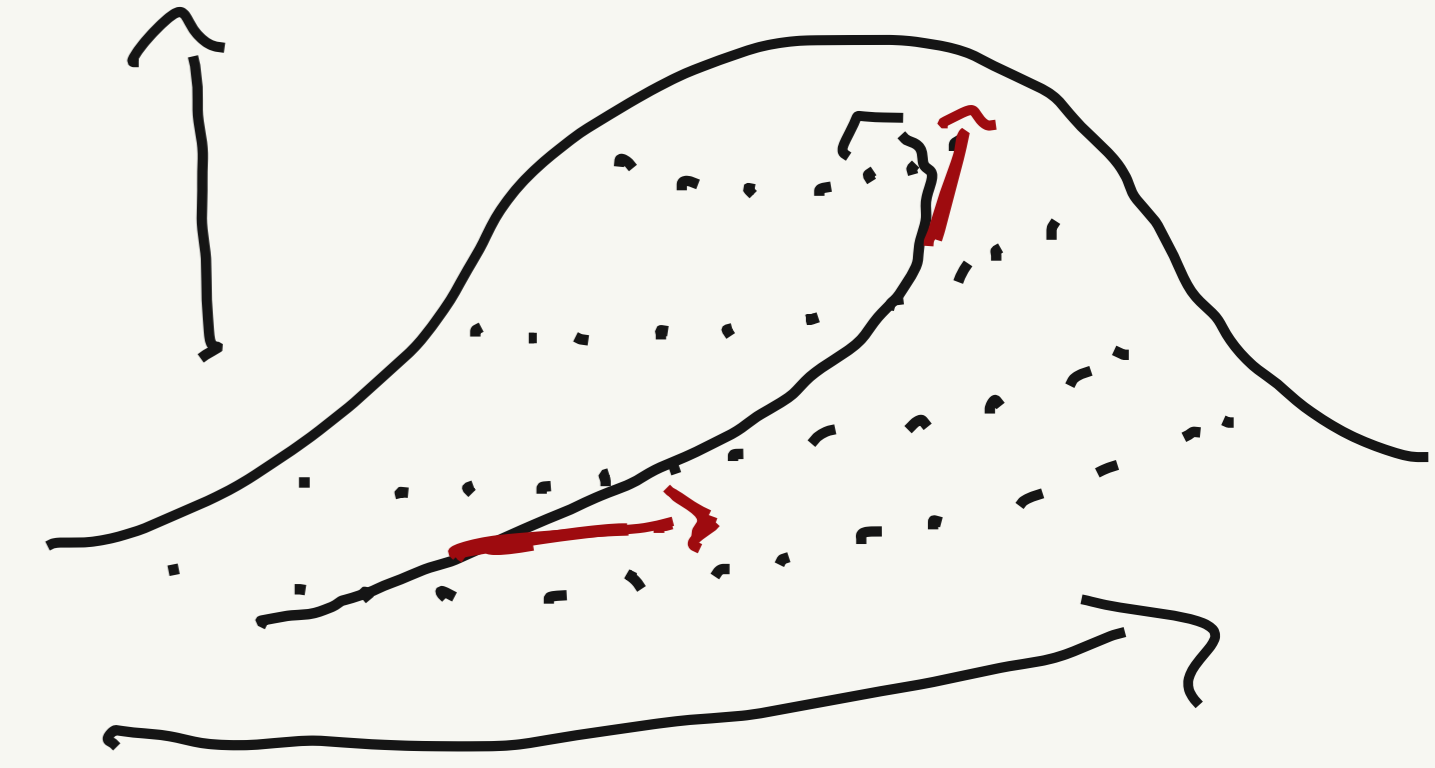
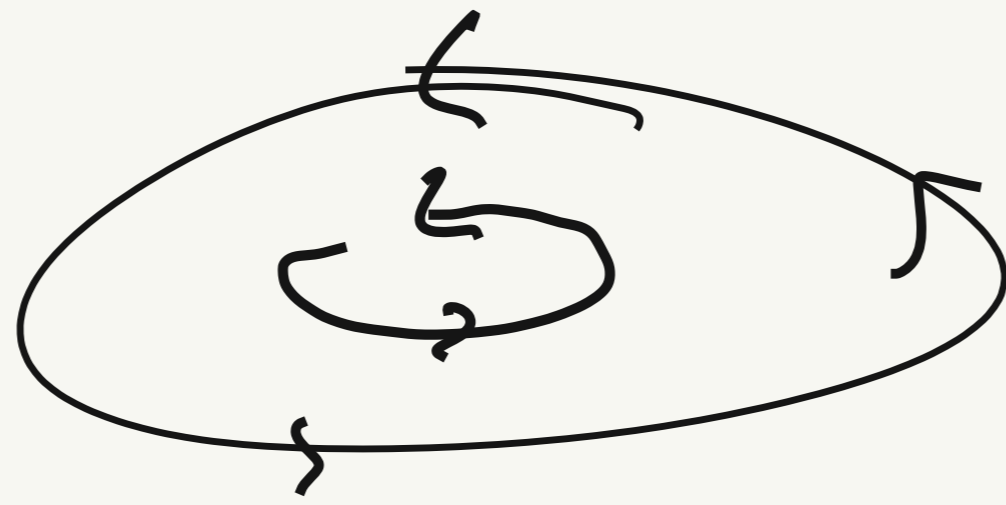
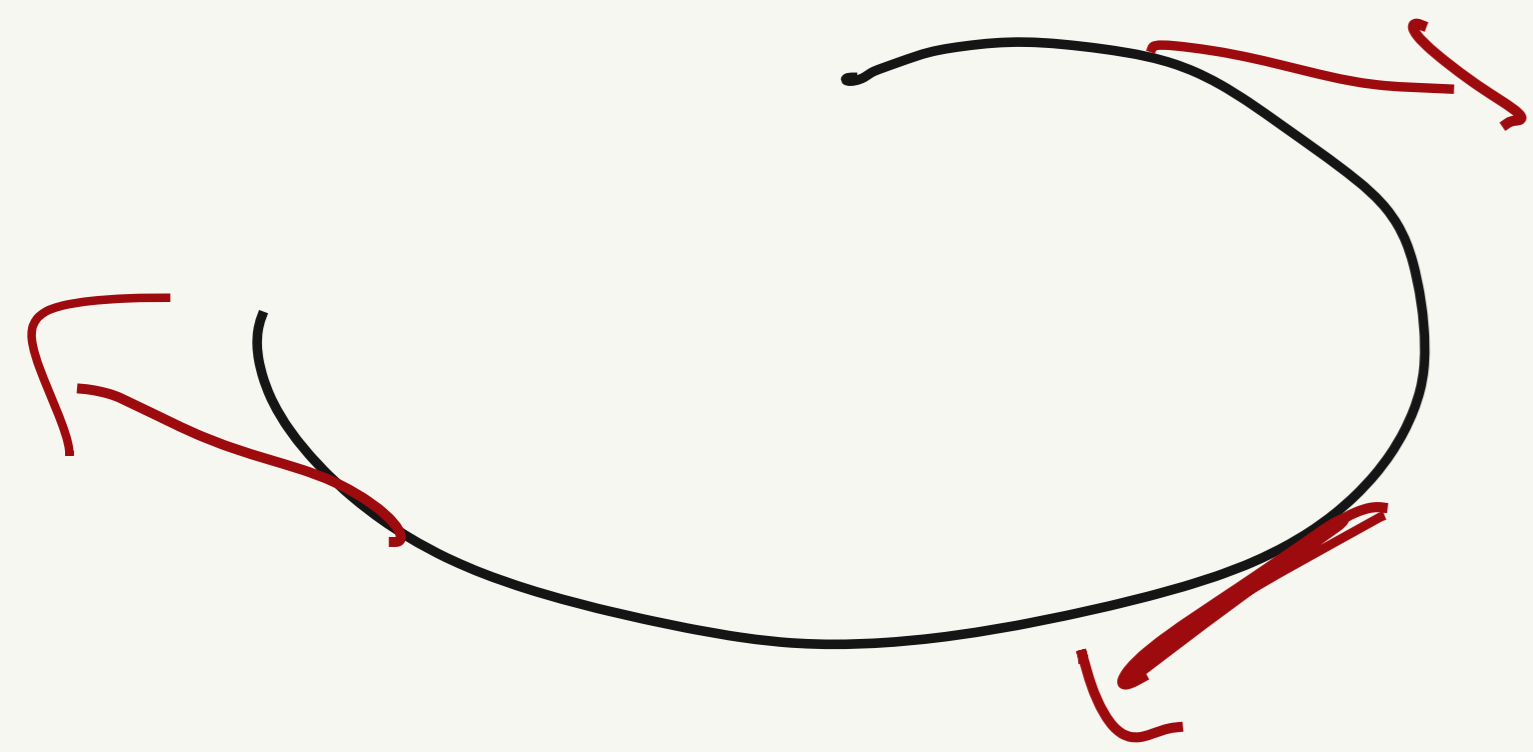
dynamics
of gauge
forced
particles

split into particle
dynamics and gauge
force

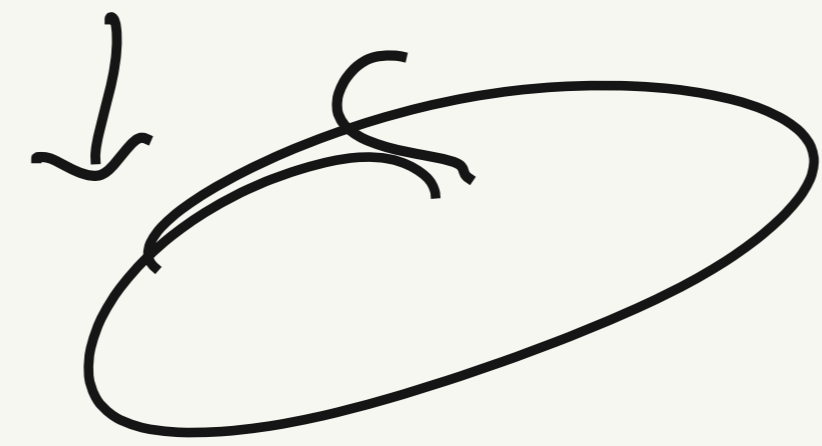
Example:

H/P

V/P



\mathbb{R}^3



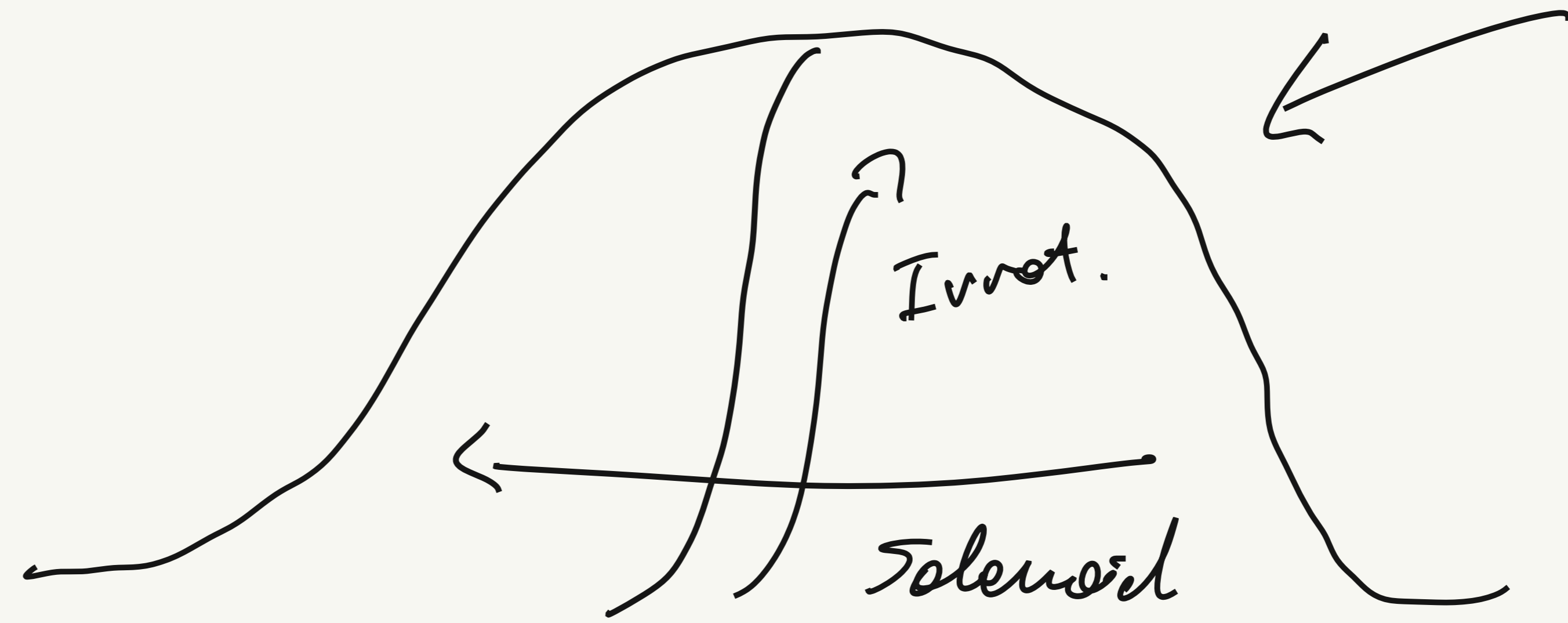
along isocont
 \Rightarrow
 max int.

up
 isocont.
 \Rightarrow
 MAP



We have
 seen this
 before





$$\psi = J(x)$$

$$\exp\{-J(x)\}$$

Taking J as $-\ln\{p(\mu)\}$,

$$\ln p = -J$$
$$\vec{p} = e^{-J}$$

this flow is $-\nabla \ln\{p(\mu)\}$.

$$-J = -J$$

Proof.

$$\partial_{x_i} p(x_i) = -\partial_{x_i} J(x_i) p(x_i)$$

$$\frac{\nabla p}{p} = \nabla \ln p$$

$$\nabla p(x_i) = -\nabla J(x_i) p(x_i)$$
$$\nabla \ln p = -\nabla J$$

What we have shown:

• $\min F[E] \Leftrightarrow \max \text{ent}$
• $\max \text{ent} \Leftrightarrow \text{potential on the state space}$

↳ Consequence:

flow towards a mode
is a sort of gauge
force.

Appendix : gauge symmetry

Frederic Schuller

Vahagn Rubakov

Leonardo A. Lezza

"Notes on Gauge theories"

2019